

EECE 380 Final Formula Sheet

Vector Formulas:

Cross Product Magnitude: $|\vec{A} \times \vec{B}| = |\vec{A}| |\vec{B}| \sin \theta$

Dot Product: $\vec{A} \cdot \vec{B} = |\vec{A}| |\vec{B}| \cos \theta$

Cartesian Coordinates: $x \rightarrow y \rightarrow z$

Cylindrical Coordinates: $r \rightarrow \phi \rightarrow z$

$$dl \phi = r d\phi$$

Spherical Coordinates: $r \rightarrow \theta \rightarrow \phi$

$$dl \theta = r d\theta, \quad dl \phi = r \sin \theta d\phi$$

Vector Projection of A on B: $\frac{\vec{A} \cdot \vec{B}}{|\vec{B}|^2} \vec{B}$

Normal Unit Vector to a Plane $ax + by + cz = cst$: $\frac{a\vec{x} + b\vec{y} + c\vec{z}}{\sqrt{a^2 + b^2 + c^2}}$

Gradient Properties: $\nabla \cdot (\nabla \times \vec{A}) = 0$

$$\nabla \times (\nabla V) = 0$$

Stoke's Theorem: $\int_S (\nabla \times \vec{A}) \cdot \vec{ds} = \oint_C \vec{A} \cdot \vec{dl}$

Divergence Theorem: $\int_V (\nabla \cdot \vec{A}) dv = \oint_S \vec{A} \cdot \vec{ds}$

Electrostatics Formulas:

Electric Field Laws:

$$A) \nabla \cdot \vec{E} = \frac{\rho_v}{\epsilon_0}$$

$$B) \nabla \times \vec{E} = 0$$

$$\text{Stoke's Theorem for } \vec{E} : \int_S (\nabla \times \vec{E}) \cdot \vec{ds} = \oint_C \vec{E} \cdot \vec{dl} = 0$$

$$\text{Divergence Theorem for } \vec{E} : \int_V (\nabla \cdot \vec{E}) dv = \oint_S \vec{E} \cdot \vec{ds} = \frac{Q}{\epsilon_0}$$

$$\text{Charge: } Q = \int_V \rho dv = \int_S \rho ds = \int_l \rho dl$$

$$\text{Field Created by a Point Charge: } \vec{E} = \frac{q}{4\pi\epsilon_0 R^2} \vec{q}_r = \frac{q}{4\pi\epsilon_0 |R_1 - R|^3} (\vec{R}_1 - \vec{R})$$

$$\text{Force between of one point charge on another: } \vec{F} = q\vec{E} = \frac{q_1 q_2}{4\pi\epsilon_0 R^2} \vec{a}_r$$

Electric field created by a:

- Line Charge: $\vec{E} = \int_l \frac{\rho_l dl}{4\pi\epsilon_0 R^2} \vec{a}_R$, Sheet: $\vec{E} = \int_s \frac{\rho_s ds}{4\pi\epsilon_0 R^2} \vec{a}_R$, Volume: $\vec{E} = \int_v \frac{\rho_v dv}{4\pi\epsilon_0 R^2} \vec{a}_R$
- Infinite Line: $\vec{E} = \frac{\rho_l}{2\pi\epsilon_0 r} \vec{a}_r$, Infinite Sheet: $\vec{E} = \frac{\rho_s}{2\epsilon_0} \vec{a}_n$, Ring: $\vec{E} = \frac{hb\rho_l}{2\epsilon_0 (b^2 + h^2)^{3/2}} \vec{a}_n$
- Disk: $\vec{E} = \frac{Ah}{2\epsilon_0} \left[\frac{-r}{\sqrt{r^2 + h^2}} + \ln \left(\frac{r + \sqrt{r^2 + h^2}}{h} \right) \right] \vec{a}_z$ (where A is the area)

$$\text{Electric field at the center of a semi-circle in the x-y plane situated above the x-axis: } \vec{E} = -\frac{\rho_l}{2\pi\epsilon_0 b} \vec{a}_y$$

$$\text{Gauss Law: } \nabla \cdot \vec{E} = \frac{\rho_v}{\epsilon_0} \rightarrow \oint_S \vec{E} \cdot \vec{ds} = \frac{Q}{\epsilon_0}$$

Electric Potential:

$$V = \frac{W}{q} = -\int \vec{E} \cdot \vec{dl}, \quad V_A - V_B = -\int_A^B \vec{E} \cdot \vec{dl}$$

$$E = -\nabla V$$

Electric Potential of a Point Charge: $V = \frac{q}{4\pi\epsilon_0 R} = \frac{q}{4\pi\epsilon_0 |\vec{r} - \vec{r}'|}$

Electric Potential for:

Line Charge: $V = \int_{\infty}^A \frac{\rho_l dl}{4\pi\epsilon_0 |\vec{R} - \vec{R}'|}$, Sheet: $V = \int_{\infty}^A \frac{\rho_s ds}{4\pi\epsilon_0 |\vec{R} - \vec{R}'|}$, Volume: $V = \int_{\infty}^A \frac{\rho_v dv}{4\pi\epsilon_0 |\vec{R} - \vec{R}'|}$

Electric Potential at a point z above a charged disk: $V = \frac{\rho_s}{2\epsilon_0} (\sqrt{r^2 + z^2} - |z|)$

Electric Dipole: $V_p = \frac{q}{4\pi\epsilon_0 R^+} - \frac{q}{4\pi\epsilon_0 R^-} \approx \frac{qd \cos \theta}{4\pi\epsilon_0 R^2} = \frac{p}{4\pi\epsilon_0 R^2} \vec{a}_r$ where $\vec{p} = q\vec{d}$

Poisson Equation: $\nabla^2 V = \frac{-\rho_v}{\epsilon_0}$

Dielectrics: $E = RJ \rightarrow J = \sigma E$

Density/Intensity Relation: $\vec{D} = \epsilon \vec{E}$

Electric Susceptibility: $\nabla \cdot \vec{E} = \frac{\rho_v + \rho_{pv}}{\epsilon_0}$ $\nabla \cdot \vec{P} = -\rho_{pv}$ $\epsilon_r = 1 + \lambda_e$
 $\vec{P} = \epsilon_0 \lambda_e \vec{E}$ $\vec{D} = \epsilon_0 \vec{E} + \vec{P} = \epsilon_0 (1 + \lambda_e) \vec{E} = \epsilon_0 \epsilon_r \vec{E} = \epsilon \vec{E}$

Boundary Conditions: $E_{2t} = E_{1t}$, $D_{1n} - D_{2n} = \rho_s$ or $\vec{a}_{21} \times (\vec{D}_1 - \vec{D}_2) = \rho_s$

$\frac{\tan \alpha_2}{\tan \alpha_1} = \frac{\epsilon_2}{\epsilon_1}$ (for a 0 charge density boundary),

Capacitance: $C = \frac{Q}{V}$

- Parallel Plate Capacitor: $C = \frac{\epsilon S}{d}$
- Toroidal Capacitor: $C = \frac{2\pi\epsilon L}{\ln\left(\frac{b}{a}\right)}$
- Spherical Capacitor: $C = \frac{4\pi\epsilon}{\frac{1}{R_i} - \frac{1}{R_o}}$

Electric Energy: $W_e = \frac{1}{2} qV = \frac{1}{2} \int_V \vec{D} \cdot \vec{E} dv = \frac{1}{2} \int_V \epsilon |\vec{E}|^2 dv$

Magnetostatics Formulas:

Magnetic Field Laws:

$$A) \nabla \cdot \vec{B} = 0$$

$$B) \nabla \times \vec{B} = \mu_0 \vec{J}$$

$$\text{Stoke's Theorem for } \vec{B} : \int_S (\nabla \times \vec{B}) \cdot \vec{ds} = \oint_C \vec{B} \cdot \vec{dl} = \mu_0 I$$

$$\text{Divergence Theorem for } \vec{B} : \int_V (\nabla \cdot \vec{B}) dv = \oint_S \vec{B} \cdot \vec{ds} = 0$$

Magnetic Field for:

- Infinite Line Current: $\vec{B} = \frac{\mu_0 I}{2\pi r} \vec{a}_\phi$
- The Inside of an Infinite Line Current: $\vec{B} = \frac{\mu_0 I_{\text{cond}} r}{2\pi b^2} \vec{a}_\phi$
- Toroidal Coil: $\vec{B} = \frac{\mu_0 NI}{2\pi r} \vec{a}_\phi$ (This value is 0 inside and outside the core)
- Finite Line Current: $\vec{B} = \frac{\mu_0 IL}{2\pi r \sqrt{L^2 + r^2}} \vec{a}_\phi$ (L is half the length of the line)
- Ring Current: $\vec{B} = \frac{\mu_0 I b^2}{2(b^2 + z^2)^{3/2}} \vec{a}_z$

$$\text{Poisson Equation: } \nabla^2 A = -\mu_0 J$$

$$\text{Vector Potential: } \vec{A} = \frac{\mu_0 I}{4\pi} \oint_C \frac{\vec{dl}}{R}$$

$$\text{Magnetic Field/Flux Density: } \vec{B} = \nabla \times \vec{A} = \frac{\mu_0 I}{4\pi} \oint_C \frac{\vec{dl} \times \vec{R}}{R^3} \text{ (Biot-Savart Law)}$$

$$\text{Magnetic Flux: } \Phi = \int_S \vec{B} \cdot \vec{ds} = \oint_C \vec{A} \cdot \vec{dl}$$

$$\text{Force: } \vec{F} = I \oint \vec{dl} \times \vec{B} \text{ or } \vec{F} = q\vec{v} \times \vec{B}$$

$$\text{Magnetic Dipole Moment: } \vec{m} = IS\vec{a}_n$$

Torque:

$$\vec{T} = \vec{d} \times \vec{F} = \vec{r} \times \vec{F}_1 + (-\vec{r}) \times \vec{F}_2$$

$$\vec{T} = \vec{m} \times \vec{B}$$

Magnetic Field Intensity of a Sheet: $\vec{H} = \frac{1}{2} \vec{J} \times \vec{a}_n$

Ampere's Law: $\oint_C \vec{H} \cdot d\vec{l} = I$

Density/Intensity Relation: $\vec{B} = \mu \vec{H}$

Magnetic Susceptibility: $\nabla \times \left(\frac{\vec{B}}{\mu_0} - \vec{M} \right) = \vec{J}$ $\vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M}$ $\mu_r = 1 + \lambda_m$
 $\vec{M} = \lambda_m \vec{H}$ $\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M} = \mu_0 (1 + \lambda_m) \vec{H} = \mu_0 \mu_r \vec{H} = \mu \vec{H}$

Boundary Conditions: $B_{1n} = B_{2n}$, $H_{1t} - H_{2t} = J_s$ or $\vec{a}_{21} \times (\vec{H}_1 - \vec{H}_2) = J_s$

$\frac{\tan \alpha_2}{\tan \alpha_1} = \frac{\mu_2}{\mu_1}$ (for a 0 charge density boundary)

Flux Linkage: $\psi = N \Phi$

Inductance: $L = \frac{\psi}{I}$

- Toroid: $L = \frac{\mu N^2 h}{2\pi} \ln\left(\frac{b}{a}\right)$
- Coaxial Transmission Line: $L = \frac{\mu I \ell}{2\pi} \ln\left(\frac{b}{a}\right)$
- Solenoid: $L = \frac{\mu N^2 S}{\ell}$

Magnetic Energy: $W_m = \frac{1}{2} qV = \frac{1}{2} \int_V \vec{B} \cdot \vec{H} dv = \frac{1}{2} \int_V \mu |\vec{H}|^2 dv$

Time Varying Fields:

$$\text{Faraday's Law: } \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \rightarrow V = -\frac{Nd\Phi}{dt} \Leftrightarrow V = \int (\vec{u} \times \vec{B}) \cdot d\vec{l} \quad (\text{Latter used for moving circuits})$$

Note: $\vec{u} \times \vec{B}$ always points in the direction of the induced current; integral is set up opposite to direction of current.

$$\text{Maxwell's Contribution: } \nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \frac{\partial \vec{D}}{\partial t}$$

$$\text{Moving Circuit with Time Varying Field: } V \simeq E = -\int_s \frac{\partial \vec{B}}{\partial t} + \oint_c (\vec{u} \times \vec{B}) \cdot d\vec{l}, \quad \vec{u} = \vec{\omega} R, \quad \alpha = \omega t$$

$$\text{Displacement and Conduction Current: } \nabla \times \vec{H} = \vec{J}_c + \vec{J}_d = \vec{J}_c + \frac{\partial \vec{D}}{\partial t} \quad \text{and} \quad \frac{J_c}{J_d} = \frac{\sigma}{\omega \epsilon}$$

$$\text{Poisson Equation: } \nabla^2 A = -\mu \vec{J} + \mu \epsilon \frac{\partial^2 \phi}{\partial t^2}$$

Plane Wave Propagation:

$$\text{Wave Equation: } A \cos\left(\frac{2\pi t}{T} - \frac{2\pi x}{\lambda} + \phi_0\right)$$

$$\text{Relations: } f = \frac{1}{T}, \quad \omega = 2\pi f = \frac{2\pi}{T}, \quad u_p = \frac{\lambda}{T} = \lambda f, \quad B = \frac{2\pi}{\lambda}$$

T: Timelength/Period (s)
 λ : Spatial/Wave length (m)
f: Frequency (Hz)
 ω : Angular Frequency (rad/s)
B: Wave Number or Phase Constant (rad/m)

Time Harmonic Fields:

$$\nabla \cdot \vec{E} = \frac{\tilde{\rho}_v}{\epsilon}$$

$$\nabla \times \vec{E} = -j \omega \vec{B} = -j \omega \mu \vec{H}$$

$$\nabla \cdot \vec{H} = \nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{H} = \vec{J} + j \omega \vec{D} = \vec{J} + j \omega \epsilon \vec{E}$$

$$\text{Complex Permeability: } \nabla \times \vec{H} = j \omega \left(\epsilon - j \frac{\sigma}{\omega} \right) \vec{E} = j \omega \epsilon_c \vec{E}, \quad \epsilon_c = \left(\epsilon - \frac{j \sigma}{\omega} \right), \quad \epsilon_{real} = \epsilon, \quad \epsilon_{imaginary} = \frac{\sigma}{\omega}$$

Maxwell's Equations in a Charge-Free Medium:

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \times \vec{E} = -j \omega \vec{B} = -j \omega \mu \vec{H}$$

$$\nabla \cdot \vec{H} = \nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{H} = j \omega \epsilon_c \vec{E}$$

$$\nabla^2 \tilde{E} + \omega^2 \mu \epsilon_C \tilde{E} = 0$$

$$\nabla^2 \tilde{E} - \gamma^2 \tilde{E} = 0$$

$$\nabla^2 \tilde{E} + k^2 \tilde{E} = 0$$

γ : Propagation Constant

k: Wavenumber

$$\gamma^2 = -\omega^2 \mu \epsilon_C, \quad k^2 = \omega^2 \mu \epsilon_C$$

$$\gamma = j \omega \sqrt{\mu \epsilon_C}, \quad k = \omega \sqrt{\mu \epsilon_C}$$

Solving:

The wave will not have a component in the direction of propagation:

$$\tilde{E}_x(z) = E_{x_0}^+ e^{-jkz} + E_{x_0}^- e^{+jkz}$$

$$\tilde{E}_y(z) = E_{y_0}^+ e^{-jkz} + E_{y_0}^- e^{+jkz}$$

$$\tilde{E}_z(z) = 0$$

Assuming no y component and only positive E_x component:

$$\tilde{H}_x = 0, \tilde{H}_z = 0$$

$$\text{Intrinsic Impedance: } \eta = \frac{\omega \mu}{k} = \sqrt{\frac{\mu}{\epsilon}}$$

$$\tilde{H}_y = -\frac{1}{j\omega \mu} \frac{\partial \tilde{E}_x(z)}{\partial z} = -\frac{1}{j\omega \mu} E_{x_0}^+ (-jk) e^{-jkz} = \tilde{H}_{y_0}^+ e^{-jkz} \quad \text{where } \tilde{H}_{y_0}^+ = \frac{k}{\omega \mu} E_{x_0}^+ = \frac{E_{x_0}^+}{\eta}$$

Wave Relations in Electromagnetics:

$$u_p = \frac{\omega}{k} = \frac{\omega}{\omega \sqrt{\mu \epsilon}} = \frac{1}{\sqrt{\mu \epsilon}}, \quad \text{In free space: } u_p = \frac{1}{\sqrt{\mu_0 \epsilon_0}} = c = 3 \times 10^8$$

Relation Between E,H, and Direction of Propagation:

$$\tilde{H} = \frac{1}{\eta} \vec{a}_k \times \tilde{E} \quad \text{or} \quad \tilde{E} = -\eta \vec{a}_k \times \tilde{H}$$

In General:

$$\begin{aligned} \tilde{E} &= \vec{a}_x E_x^+(z) + \vec{a}_y E_y^+(z) \\ \tilde{H} &= \vec{a}_x H_x^+(z) + \vec{a}_y H_y^+(z) \end{aligned} \quad \text{so:} \quad H_x^+(z) = \frac{-E_y^+(z)}{\eta} \quad \text{and} \quad H_y^+(z) = \frac{-E_x^+(z)}{\eta}$$

Wave Polarization:

$$\text{Let } \tilde{E}(z) = \vec{a}_x \tilde{E}_x + \vec{a}_y \tilde{E}_y \quad \text{with} \quad \begin{aligned} \tilde{E}_x(z) &= E_{x_0}^+ e^{-jkz} = |E_{x_0}^+| e^{-jkz} e^{+j\phi_x} \\ \tilde{E}_y(z) &= E_{y_0}^+ e^{-jkz} = |E_{y_0}^+| e^{-jkz} e^{+j\phi_y} \end{aligned}$$

$$\text{With } E_x \text{ as reference: } E_{x_0}^+ = A_x \quad \text{and} \quad E_{y_0}^+ = |E_{y_0}^+| e^{j\delta} = A_y e^{j\delta}$$

$$E(z, t) = \vec{a}_x A_x \cos(\omega t - kz) + \vec{a}_y A_y \cos(\omega t - kz + \delta)$$

$$|E(z, t)| = \sqrt{(A_x \cos(\omega t - kz))^2 + (A_y \cos(\omega t - kz + \delta))^2}$$

$$\Psi(z, t) = \tan^{-1} \left(\frac{A_y \cos(\omega t - kz + \delta)}{A_x \cos(\omega t - kz)} \right) \text{ (usually } z \text{ is taken equal to } 0)$$

- For Linear Polarization (z taken equal to 0):

$$\delta = 0 \text{ or } \delta = \pi$$

$$\text{If } \delta = 0: \Psi(0, t) = \tan^{-1} \left(\frac{A_y \cos(\omega t)}{A_x \cos(\omega t)} \right) = \frac{A_y}{A_x} \text{ (constant)}$$

$$\text{If } \delta = \pi: \Psi(0, t) = \tan^{-1} \left(\frac{A_y \cos(\omega t + \pi)}{A_x \cos(\omega t)} \right) = -\frac{A_y}{A_x} \text{ (constant)}$$

$$\text{In both cases: } |E(0, t)| = \cos t(\omega) \sqrt{A_x^2 + A_y^2}$$

- For Circular Polarization (z taken equal to 0 and $A_x = A_y = A$):

$$\delta = \frac{\pi}{2} \text{ or } \delta = -\frac{\pi}{2}$$

$$\text{If } \delta = \frac{\pi}{2} \text{ Left Hand Circular Polarization: } \Psi(0, t) = \tan^{-1} \left(-\frac{A \sin(\omega t)}{A \cos(\omega t)} \right) = -\omega t$$

$$\text{If } \delta = -\frac{\pi}{2} \text{ Right Hand Circular Polarization: } \Psi(0, t) = \tan^{-1} \left(\frac{A \sin(\omega t)}{A \cos(\omega t)} \right) = \omega t$$

$$\text{In both cases: } |E(0, t)| = A \sqrt{\cos^2(\omega t) + \sin^2(\omega t \pm \frac{\pi}{2})} = A \text{ (constant)}$$

- For Elliptical Polarization:

δ is neither of the previous angles, both the magnitude and angle are not constant.

If $\sin \delta > 0$ left hand elliptical polarization.

If $\sin \delta < 0$ right hand elliptical polarization.

Constants:

$$\varepsilon_0 = \frac{10^{-9}}{36\pi} (F / m)$$

$$\mu_0 = 4\pi \times 10^{-7} (H / m)$$

$$\eta = \sqrt{\frac{\mu_0}{\varepsilon_0}} = 120\pi (\Omega) \text{ (in free space)}$$

$$c = 3 \times 10^8 (m / s)$$